REPORT

EXOPLANETS

A multiple planet system of super-Earths orbiting the brightest red dwarf star GJ 887

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The closet exoplanets to the Sun provide opportunities for detailed characterization of planets outside the Solar System. We report the discovery, using radial velocity measurements, of a compact multiplanet system of super-Earth exoplanets orbiting the nearby red dwarf star GJ 887. The planets have orbital periods of 9.3 and 21.8 days. Assuming an Earth-like albedo, the equilibrium temperature of the 21.8-day planet is ~350 kelvin. This temperature is within, but close to the inner edge of, the liquid-water habitable zone. We also detect an unconfirmed signal with a period of ~50 days, which could correspond to a third super-Earth in a more temperate orbit. Our observations show that GJ 887 has photometric variability below 500 parts per million, which is unusually quiet for a red dwarf.

t visible wavelengths, GJ 887 (HD 217987) is the brightest red dwarf in the sky and, at a distance of 3.29 pc, is the 12th closest star system to the Sun. GJ 887 is the most massive red dwarf within 6 pc of the Sun, which is close enough for a direct stellar radius measurement using interferometry (1). GJ 887's stellar parameters are listed in Table 1. Red dwarfs are amenable to radial velocity (RV) searches for temperate Earth-mass exoplanets: their low luminosity means that temperate planets have short orbital periods, and

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their low stellar mass implies that Earth-mass planets can impart a reflex RV detectable with current instrumentation. Although the transit method of planet discovery efficiently detects planets because many stars can be simultaneously monitored, it will detect only planets that pass through the line of sight between Earth and the host star. Consequently, only 1 to 2% of habitable-zone planets (i.e., those with surfaces that can support liquid water) are detectable with the transit method. The RV method is necessary to achieve a complete census of the planets orbiting our closest stellar neighbors, especially red dwarfs.

We monitored GJ 887 as part of the Red Dots #2 project. Nightly observations were taken with the High Accuracy Radial Velocity Planet Searcher (HARPS) (2) for 3 months. We also obtained contemporaneous photometric observations (3). Regular nightly sampling combined with photometric observations mitigates against false-positive exoplanet detections caused by intrinsic stellar variability and other sources of correlated noise. We supplement our data with >200 archival observations from HARPS, the Planet Finder Spectrograph (PFS) (4), the High-Resolution Echelle Spectrometer (HIRES) (5), and the University College London Echelle Spectrograph (UCLES) (6), spanning nearly 20 years (3). We used photometry from various ground-based observatories and the Transiting Exoplanet Survey Satellite (TESS) spacecraft (7). Tables S3 and S4 list all the data we used.

We searched for a candidate planet by adding a (circular) Keplerian orbit test signal to a base model and measuring the improvement in the logarithm of the likelihood statistic (*3*). The base model is composed of an offset and an instrumental jitter added to the measurement uncertainties for each dataset. We used this to generate log-likelihood periodograms for both the RV and photometric data and then searched for signals by plotting the increase in the log-likelihood statistic against test period (Fig. 1). The highest peaks were evaluated for statistical significance (*8*, *9*). We recursively added additional planet test signals, adjusting all the parameters to maximize the likelihood

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Fig. 1. Periodograms of the RV data. (A) The log-likelihood (Δ In L) periodogram obtained separately for all RV data before 2018 (brown) and the Red Dots #2 campaign (red). (**B**) the same search for a first signal when combining all the RV observations together. The vertical green lines indicate our derived model periods for planets b and c and the candidate planet d. The horizontal dashed lines in both panels indicate the false alarm probability (FAP) values.

of all planet signals and the parameters of the base model. We continued this iterative process until no signals below a threshold of 0.1% false-alarm probability were found in the time series (fig. S1). We detected periodic signals at 9.3, 21.8, and 50.7 days (Fig. 1B) and verified them using several independent fitting procedures and algorithm implementations (*3*). The regular sampling of the Red Dots #2 dataset helped to disentangle the signals under investigation (Fig. 1A).

Stellar magnetic activity can induce an asymmetric distortion of the spectral lines, shifting the measured line center and consequently inducing an apparent RV shift, which may appear as a false-positive exoplanet at the stellar rotation period (10). The rotation period of GJ 887 is unknown, so we searched for periodicities in the photometric data (3). The archival data from 2002-2004 show an ~200-day period, but this period was undetectable in the 2018 quasi-simultaneous photometric observations because the time span is too short. Our analysis of the TESS photometry shows very low intrinsic variability, with a semi-amplitude of 240 parts per million. It is unclear whether this is caused by systematic effects, known to affect the TESS observations, but we use this value as an upper limit to the intrinsic variability of GJ 887. The TESS variability can be explained by one starspot, or a group of starspots, with a total diameter of 0.3% of the stellar surface, indicating that GJ 887 is slowly rotating with very few surface brightness inhomogeneities (*11*). GJ 887 is less magnetically active than most stars with the same effective temperature, as demonstrated by: the very low starspot coverage; low photometric variability; the activity metric derived from stellar Ca II H and K lines, $\log(R'_{\rm HK}) = -4.805$ (*12*); and the very low H α activity (*13*).

The RV signals are detected in the Red Dots # 2 HARPS data alone (Fig. 1A), so we investigated additional spectral signatures of stellar magnetic activity using this dataset. We extracted a time series of the flux in the cores of the Na I D, H α , and H β lines; and calculated the S-index, which is the ratio of flux in the cores in the Ca II H and K lines compared with the continuum (3). The S-index and Na D lines both show a weak signal at about 55 days, whereas the H α and H β lines show a weak signal at 38 days (fig. S5). These differing periodicities could reflect time scales of various stellar activity processes on the star, and despite being low in amplitude, they make a planetary origin for RV signals in the 30-to-60-day domain less certain. None of these activity periodicities are close to the RV signals at 9.3 and 21.8 days, but they do make the RV signal at 50.7 days questionable.

Correlated noise, for example, that caused by stellar activity, can be assessed using the covariances between observations. To further test the planetary origin of the detected RV signals, we fitted maximum likelihood model functions using two planet models with and without Gaussian processes (GPs) (3). All of the models including GPs improved the fit to the data compared with those without, and the amplitude of the signals with periods of 9.3 and 21.8 days remained unchanged within their 1σ uncertainty. The modeling of the correlated noise using GPs therefore does not affect these two signals. However, the significance of the third signal drops substantially when including a GP in the model, casting further doubts on its Keplerian nature. Table S4 lists the derived values and statistics from these models.

We conclude that the two signals with orbital periods of 9.3 and 21.8 days correspond to two exoplanets, planets b and c, respectively. The minimum masses of these planets (m_p) are 4.2 ± 0.6 and 7.6 ± 1.2 Earth masses (M_{\oplus}) , which makes them two super-Earth exoplanets with orbital semimajor axes (a_p) of 0.068 and 0.120 astronomical units (au). The inner



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planet has an orbital eccentricity consistent with zero (Fig. 2), but the outer planet is more likely to have a small nonzero eccentricity (Fig. 3). We regard the third signal at ~50 days (compare Fig. 2) as dubious and likely related to stellar activity. The fits to our two-planet model and the two-planet + third signal model are shown in Fig. 2.

The long-term dynamical stability of the orbits can also be used to test the physical plausibility of the system and investigate whether it contains unusual configurations such as dynamical resonances. We performed a dynamical stability study (3) using the software package MERCURY6 (14). We find that all two-planet solutions are stable even if eccentricities are allowed to vary. The ratio of periods of these two planets is close to 7:3, but the simulations do not support the presence of a dynamical resonance, as there is an absence of oscillating orbital alignment variations (15). However, we find that the system must be in a dynamically active state, driving oscillatory changes in the eccentricities of both planets. These interactions produce very regular variations, supporting the hypothesis that the two-planet configuration is dynamically stable on very long time scales. For a putative system with three planets, only ~25% of the 1000 best-fitting models would

Table 1. Stellar parameters for GJ 887 and parameters for planets b and c. The top half of the table lists parameters for GJ 887: the parallax in milliarcseconds (mas), distance in parsecs (pc), *V*- and *G*-band magnitudes, stellar mass in solar masses (M_{\odot}), metallicity relative to the Sun [Fe/H], luminosity and radius in solar units, projected rotational velocity (vsini), and surface gravity log *g*. The stellar mass was computed using a mass-radius relation (*26*). The bottom half of the table lists parameters for planets b and c: S_{eff} is the incident flux from GJ 887 in units of the incident flux on Earth from the Sun, and T_{eq} is the equilibrium temperature of the planet.

Parameter	Value	Reference
Spectral type	M1V	(27)
Parallax (mas)	304.2190 ± 0.0451	(28)
Distance (pc)	3.2871 ± 0.0005	
Magnitude	V = 7.34, G = 6.522	(28)
Mass (M_{\odot})	0.489 ± 0.05	
[Fe/H]	-0.06 ± 0.08	(27)
T _{eff} (K)	3688 ± 86	(27)
Luminosity (L_{\odot})	0.0368 ± 0.004	(27)
Radius (R_{\odot})	0.4712 ± 0.086	(1, 29, 30)
vsini (km s ⁻¹)	2.5 ± 1.0	(31)
logR' _{HK} mean	-4.805 ± 0.023	(12)
log (age/years)	9.46 ± 0.58	(27)
log g	4.78	(32)
	GJ 887 b	GJ 887 c
$K_{\rm p} ({\rm m \ s^{-1}})$	2.1 ^{+0.3}	2.8 ± 0.4
<i>P</i> _p (d)	9.262 ± 0.001	21.789 ^{+0.004} -0.005
$\overline{m_{\rm p}}(M_{\oplus})$	4.2 ± 0.6	7.6 ± 1.2
a _p (au)	0.068 ± 0.002	0.120 ± 0.004
$S_{eff,p} (S_{eff,\oplus})$	7.95 ± 0.2	2.56 ± 0.2
T _{eq} (K)	468	352

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be dynamically stable over 10^5 years, but this is mostly due to the poorly constrained eccentricities. Given that the observations only provide upper limits on the eccentricities, we investigated orbits that are assumed to be initially circular (initial zero eccentricities). Even in this three-planet case, >99% of configurations were found to be stable, meaning that the presence of a third planet cannot be ruled out using dynamic stability considerations.

The separations between the planets, in units of their spheres of gravitational influence (Hill radii), are ~19.1 for planets b and c and ~17.2 for planets c and d (if planet d is real and has a mass of 8.3 M_{\oplus}); these values are consistent with the system having undergone dynamical relaxation (16). Dynamical relaxation in systems of super-Earths results in ~80% of planets having orbital eccentricities $e_{\rm p} \leq 0.1$, with the remaining 20% having $e_{\rm p} \leq$ 0.3 (17). We examined the expected tidal evolution of GJ 887 b using analytical methods (18, 19), finding that the tidal circularization time scale of GJ 887 b is a few billion years for an assumed tidal dissipation parameter $Q'_{\rm p} =$ 1000. This is consistent with our observation that GJ 887 b's orbit is almost circular.

The multiplanet super-Earth system around GJ 887 is consistent with recent planet formation models (20, 21). These models typically form chains of multiple planets trapped in mean-motion resonances that then migrate into orbits close to the central star. Depending on where the initial planets formed in the protoplanetary disc, they could have accreted either large amounts of water ice or just dry rocky silicates. As such, the planets may be either water-rich or water-poor. At the end of the gas disc lifetime, the resonant chains of planets can remain stable, yielding systems similar to the seven-planet TRAPPIST-1 planetary system (22), or they can become unstable, leading to collisions between planets and thus a nonresonant configuration (20). The GJ 887 planetary system appears to be more consistent with the latter scenario of unstable evolution. The presence of dynamical resonances can be very sensitive to the existence or absence of additional planets. Consequently, if the third signal at 50.7 days is real or if additional planets exist, this may result in a more resonant system.

We used standard methods (23) to calculate the distances from GJ 887 within which planets could support liquid water on their surfaces [the star's habitable zone (HZ)] and found that it extends from ~0.19 to 0.38 au. With an a_p of 0.120 ± 0.004, GJ 887 c is closer to its host star than the HZ but near the inner edge. If the ~50-day signal is planetary in origin, it corresponds to a super-Earth in GJ 887's HZ. Assuming an albedo, α , similar to Earth's (α = 0.3), the equilibrium temperature, T_{eep} of planets b and c would be 468 and 352 K, respectively.

Their incident energy fluxes from the star (the insolation S) are 7.95 and 2.56 times the Sun's insolation on Earth. Figure 3 shows the insolation of known planets orbiting M-type red dwarfs as a function of host star apparent magnitude. GJ 887 has the brightest apparent magnitude of any known M dwarf planet host. This brightness, combined with the high photometric stability of GJ 887, exhibited in the TESS data, and the high planet-star brightness and radius ratios, make these planets potential targets for phase-resolved photometric studies, especially in emission (24). Spectrally resolved phase photometry has been shown to be sensitive to the presence of an atmosphere and molecules such as CO_2 (25).

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SUPPLEMENTARY MATERIALS

science.sciencemag.org/content/[vol]/[issue]/[page]/suppl/DC1 Materials and Methods Figs. S1 to S11 Tables S1 to S7 References (*34–70*) Data S1 9 August 2019; accepted 12 May 2020 10.1126/science.aaz0795

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